

Chapter-3 - Dynamic of a Particle. General Motion

① $\text{grad}(\phi + \psi)$

$$\nabla(\phi + \psi) = \nabla\phi + \nabla\psi$$

② $\text{div}(\vec{A} + \vec{B})$

$$\nabla \cdot (\vec{A} + \vec{B}) = \nabla \cdot \vec{A} + \nabla \cdot \vec{B}$$

③ $\text{curl}(\vec{A} + \vec{B})$

$$\nabla \times (\vec{A} + \vec{B}) = \nabla \times \vec{A} + \nabla \times \vec{B}$$

④ $\nabla \cdot (\phi \vec{A}) = (\nabla\phi) \cdot \vec{A} + \phi(\nabla \cdot \vec{A})$

⑤ $\nabla \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B})$

⑥ $\text{div grad } \phi$

$$\nabla \cdot (\nabla\phi) = \nabla^2\phi = \frac{d^2\phi}{dx^2} + \frac{d^2\phi}{dy^2} + \frac{d^2\phi}{dz^2}$$

⑦ $\text{curl grad } \phi$

$$\nabla \times (\nabla\phi) = \text{zero}$$

⑧ $\text{div curl } \vec{A}$

$$\nabla \cdot (\nabla \times \vec{A}) = \text{zero}$$

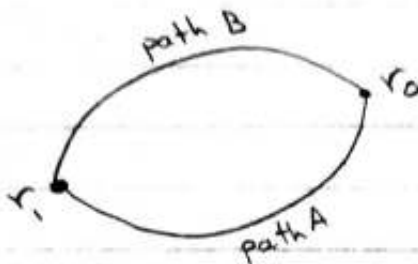
⑨ $\nabla = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z}$

'Conservation Forces''

القوة المحافضة

The condition of conservation force

- ① if the work is not depended for the path the Force is conservation and its field called conservation field.



* W_A ، شغل القوة لنقل جسم من

r_1 إلى r_0 على مسار A

$$W_A = \int_{r_1}^{r_0} \vec{F} \cdot d\vec{r}$$

* W_B ، شغل القوة لنقل جسم من

$$W_B = \int_{r_1}^{r_0} \vec{F} \cdot d\vec{r}$$

$$W_{\text{path A}} = W_{\text{path B}}$$

$$\int_{r_1}^{r_0} \vec{F} \cdot d\vec{r} = \int_{r_1}^{r_0} \vec{F} \cdot d\vec{r}$$

$$\int_{r_1}^{r_0} \vec{F} \cdot d\vec{r} - \int_{r_1}^{r_0} \vec{F} \cdot d\vec{r} = 0$$

$$\int_{r_1}^{r_0} \vec{F} \cdot d\vec{r} + \int_{r_0}^{r_1} \vec{F} \cdot d\vec{r} = 0$$

$$\oint \vec{F} \cdot d\vec{r} = 0$$

بشرط أن يكون شغل

القوة المحافضة

إذا كان لشغل القوة حول مسار مغلق يساوي صفر تكون القوة محافظة.

(3) if $\vec{F} = -\nabla V$

إذا كانت القوة تساوي سالب التدرج لجهد تكون القوة محافظة

(4) if $\nabla \times \vec{F} = 0$ then the force is Conservation

(5) if $E = K.E + P.E = \text{constant}$
then the force is conservation

ex)

$$\vec{F} = \hat{i}(y^3z^3 - 6xz^2) + \hat{j}(2xyz^3) + \hat{k}(3xy^3z^3 - 6x^2z)$$

show that if the force is conservation or no?

$$\nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ F_x & F_y & F_z \end{vmatrix}$$

$$\nabla \times \vec{F} = \hat{i} \left(\frac{\partial F_z}{\partial y} - \frac{\partial F_y}{\partial z} \right) + \hat{j} \left(\frac{\partial F_x}{\partial z} - \frac{\partial F_z}{\partial x} \right) + \hat{k} \left(\frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y} \right)$$

$$= \hat{i} (6xy^3z^3 - 6xy^3z^3) + \hat{j} [(3y^4z^3 - 12xz) - (3y^3z^3 - 12xz)] + \hat{k} (2yz^3 - 2yz^3)$$

$$= \text{Zero}$$

∴ القوة محافظة

if
ex) $V = x^2 + yx + xz^2$ Find the force field?

$$\vec{F} = -\nabla V = -\hat{i} \frac{\partial V}{\partial x} - \hat{j} \frac{\partial V}{\partial y} - \hat{k} \frac{\partial V}{\partial z}$$

$$\vec{F} = -\hat{i}(2x+y+z) - \hat{j}x - \hat{k}2x$$

ex) show that if the force is conservation or no?

$$\vec{F} = \hat{i}xy + \hat{j}xz + \hat{k}yz$$

$$\nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ xy & xz & yz \end{vmatrix}$$

$$= \hat{i}(z-x) - \hat{j}(0) + \hat{k}(z-x).$$

∴ النتيجة لا تساوي صفر. فالقوة غير محافظة.

ex) when the force field is conservation?

$$\vec{F} = \hat{i}(ax+by^2) + \hat{j}cxy$$

$$\nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ ax+by^2 & cxy & 0 \end{vmatrix} = \hat{k}(c-2b)y$$

if the force is conservation then $\nabla \times \vec{F} = 0$

$$\therefore c-2b=0 \Rightarrow c=2b$$

* Motion of a projectile in a Uniform Gravitational Field

(i) No Air Resistance

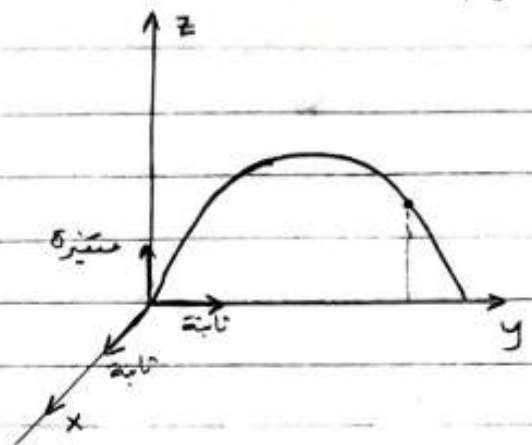
For a projectile there is only one force acting namely the force of gravity and equal the weight.

إذا كانت القوة الوحيدة المؤثرة على الجسم هي قوة الجاذبية والتي تساوي في هذه الحالة وزن الجسم، فإن الجسم يتحرك سقوطاً حراً.

the differential equation of motion:.

$$m \frac{d^2 \vec{r}}{dt^2} = -mg \hat{k}$$

$$\frac{d^2 \vec{r}}{dt^2} = -g \hat{k}$$



بما أن المعادلة لا تحتوي على الكتلة (m) فإن حركة الجسم لا تعتمد على كتلته.

the force function is clearly of the separable type and is also conservative.

From Force conservation Law?

$$\frac{1}{2} m v_0^2 = \frac{1}{2} m v^2 + m g z \quad] \times \frac{2}{m}$$

$$v^2 = v_0^2 - 2 g z$$

حركة الجسيم في أية لحظة (1) -
وهي نفس معادلة الحركة الجذبية

معادلة التفاضلية

$$\frac{d\vec{r}}{dt} = -g \hat{k} \Rightarrow \frac{d\vec{v}}{dt} = -g \hat{k}$$

$$d\vec{v} = -g \hat{k} \int dt \Rightarrow \vec{v} = -gt \hat{k} + \text{const}$$

if $t=0$, $\therefore v=v_0$, $\therefore \text{const} = v_0$

$$\boxed{\begin{aligned} \vec{v} &= v_0 - gt \hat{k} \\ v &= v_0 - gt \end{aligned}} \Rightarrow (2) \quad \text{تعبير سرعة الحركة الخطية}$$

$$\vec{v} = \frac{d\vec{r}}{dt} = v_0 - gt \hat{k}$$

$$\int d\vec{r} = \int v_0 dt - gt dt \hat{k}$$

$$\boxed{r = v_0 t - \frac{1}{2} gt^2 \hat{k}} \quad (3) \quad \text{معادلة الحركة الخطية}$$

$$\therefore x = x_0 t$$

$$y = y_0 t$$

$$z = z_0 t - \frac{1}{2} gt^2$$

Linear Air Resistance.

The differential equation is:

$$m \frac{d^2 \vec{r}}{dt^2} = -mg \hat{k} - c \vec{v}$$

$$m(\hat{i}x'' + \hat{j}y'' + \hat{k}z'') = -mg\hat{k} - c(\hat{i}x' + \hat{j}y' + \hat{k}z')$$

$$\therefore mx'' = -cx' \Rightarrow x'' = -\frac{c}{m} x'$$

$$(x'' = -\gamma x') \text{ --- (1)}$$

$$\frac{c}{m} = \gamma$$

$$\therefore my'' = -cy' \Rightarrow y'' = -\frac{c}{m} y'$$

$$(y'' = -\gamma y') \text{ --- (2)}$$

$$\therefore mz'' = -mg - cz'$$

$$(z'' = -\gamma z' - g) \text{ --- (3)}$$

to find the velocity in (x) direction

$$x'' = -\gamma x' \text{ --- (1)}$$

$$\frac{dx'}{dt} = -\gamma x'$$

$$\int_{x_0}^{x'} \frac{dx'}{x'} = -\gamma \int_0^t dt$$

$$\ln \frac{x'}{x_0} = -\gamma t$$

$$(x' = x_0 e^{-\gamma t}) \text{ --- (1')}$$

From eq. (2)

$$y'' = -\gamma y' \Rightarrow \frac{dy'}{dt} = -\gamma y' \Rightarrow \int_{y_0}^{y'} \frac{dy'}{y'} = -\gamma \int_0^t dt$$

$$\ln \frac{y'}{y_0} = -\gamma t$$

$$(y' = y_0 e^{-\gamma t}) \quad (2')$$

From eq. (3)

$$\frac{dz'}{dt} = -\gamma z' - g \Rightarrow \int_{z_0}^{z'} \frac{dz'}{-\gamma z' - g} = \int_0^t dt$$

$$\frac{1}{-\gamma} \int_{z_0}^{z'} \frac{-\gamma dz'}{-\gamma z' - g} = \int_0^t dt$$

$$\ln \frac{(-\gamma z' - g)}{(-\gamma z_0 - g)} = -\gamma t$$

$$\ln \frac{\gamma z' + g}{\gamma z_0 + g} = -\gamma t$$

$$\gamma z' + g = (\gamma z_0 + g) e^{-\gamma t}$$

$$(z' = \frac{-g}{\gamma} + (z_0 + \frac{g}{\gamma}) e^{-\gamma t}) \quad (3')$$

* to Find the position Components
From eq. (i)

$$\dot{x} = x_0 \cdot e^{-\gamma t} \Rightarrow \frac{dx}{dt} = x_0 \cdot e^{-\gamma t}$$

$$\int_{x_0}^x dx = \int_0^t x_0 \cdot e^{-\gamma t} dt \quad \left[\frac{-x}{-\gamma} \right]$$

$$x \Big|_{x_0}^x = \frac{x_0}{-\gamma} \int_0^t e^{-\gamma t} (-\gamma) dt \Rightarrow x - x_0 = \frac{-x_0}{\gamma} \left[e^{-\gamma t} \right]_0^t$$

إذا كان الموضع عند بداية في نقطة الأصل

$$x = \frac{-x_0}{\gamma} \left[e^{-\gamma t} - 1 \right]$$

$$x = \frac{x_0}{\gamma} \left[1 - e^{-\gamma t} \right] \text{ --- (1'')}$$

وبنفس الطريقة نجد (y)

$$y = \frac{y_0}{\gamma} \left(1 - e^{-\gamma t} \right) \text{ --- (2'')}$$

دبر بار (z) تأخذ المعادلة (3')

$$\dot{z} = \frac{-g}{\gamma} + \left(\frac{g}{\gamma} + z_0 \right) e^{-\gamma t} = \frac{dz}{dt}$$

$$\int_{z_0}^z dz = \frac{-g}{\gamma} \int_0^t dt + \left(\frac{g}{\gamma} + z_0 \right) \int_0^t e^{-\gamma t} dt \quad \text{نضرب بنسبة (-\gamma)}$$

$$z - z_0 = \frac{-g}{\gamma} t + \left(\frac{1}{-\gamma} \right) \left(\frac{g}{\gamma} + z_0 \right) \int_0^t e^{-\gamma t} (-\gamma) dt$$

عند نقطة الأصل $z_0 = z_{\text{zero}}$

$$z = \frac{-g}{\gamma} t - \frac{1}{\gamma} \left(\frac{g}{\gamma} + z_0 \right) \left[e^{-\gamma t} \right]_0^t$$

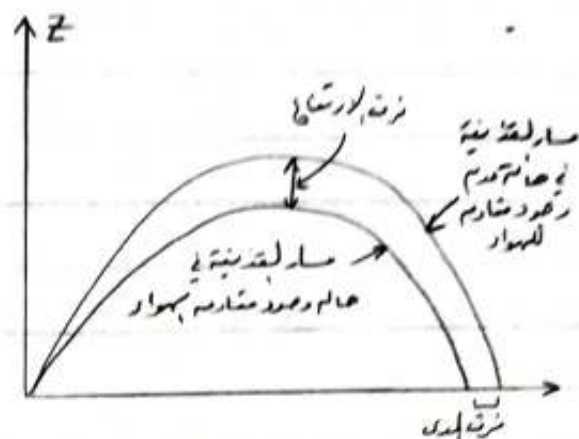
$$z = \frac{-g}{\gamma} t - \left(\frac{g}{\gamma^2} - \frac{z_0}{\gamma} \right) \left[e^{-\gamma t} - 1 \right]$$

$$z = \frac{-g}{\gamma} t + \left(\frac{g}{\gamma^2} + \frac{z_0}{\gamma} \right) \left(1 - e^{-\gamma t} \right) \text{ --- (3'')}$$

وبعد بهم المعادلات x, y, z في نفس المكان

$$\vec{r} = \hat{i}x + \hat{j}y + \hat{k}z$$

$$\vec{r} = \left(\frac{\vec{U}_0}{\gamma} + \frac{\vec{k}g}{\gamma^2} \right) (1 - e^{-\gamma t}) - \hat{k} \frac{gt}{\gamma}$$



في حالة وجود مقاومة هواء صغيرة ($\gamma t \ll 1$)
مستفيدة المعادلات

$$e^u = 1 + u + \frac{u^2}{2!} + \frac{u^3}{3!} + \frac{u^4}{4!} + \dots$$

مكثف الناتج

$$(1 - e^{-\gamma t}) = \gamma t - \frac{\gamma^2 t^2}{2!} + \frac{\gamma^3 t^3}{3!}$$

وبعد بوضعها في المعادلات (2)

$$\vec{r} = \vec{U}_0 t - \frac{1}{2} g t^2 \hat{k} - \Delta \vec{r}$$

حيث $(\Delta \vec{r})$ يمثل جميع الحدود المتبقية وهو عامل إهمال في مسار القذيفة

$$\Delta \vec{r} = \gamma \left[\vec{U}_0 \left(\frac{t^2}{2!} - \frac{\gamma t^3}{3!} + \dots \right) + \frac{g}{\gamma} \left(\frac{t^3}{3!} - \frac{\gamma t^4}{4!} + \dots \right) \right]$$

ملاحظة / فإني المعادلة لا تنطبق التي تمثل $(\vec{r})$ بوجود مقاومة هواء بالمعادلة
لباقية (بدون وجود مقاومة هواء)

$$\vec{r} = \vec{U}_0 t - \frac{1}{2} g t^2 \hat{k}$$